Nonlinear optics: a first look

J. E. Sipe Department of Physics University of Toronto



"Optical alchemy"





Cambridge Studies in Modern Optics 9 The Elements of Nonlinear Optics



Nonlinear Optical Systems

Luigi Lugiato, Franco Prati and Massimo Brambilla

Fundamentals of Nonlinear Optics

SECOND EDITION



Peter E. Powers Joseph W. Haus



Susceptibilities $\chi^{(2)}$ effects $\chi^{(3)}$ effects Quantum nonlinear optics Nonlinear optics and electronics Forbidden processes Susceptibilities $\chi^{(2)}$ effects $\chi^{(3)}$ effects Quantum nonlinear optics Nonlinear optics and electronics Forbidden processes

$$P(t) = \chi E(t)$$

$$P(t) = \chi^{(1)} E(t)$$

$$P(t) = \chi^{(1)}E(t) + \chi^{(2)}E^2(t) + \chi^{(3)}E^3(t) + \dots$$

$P(t) = \chi^{(1)}E(t) + \chi^{(2)}E^2(t) + \chi^{(3)}E^3(t) + \dots$



$P(\mathbf{r},t) = \chi^{(1)}(\mathbf{r})E(\mathbf{r},t) + \chi^{(2)}(\mathbf{r})E^{2}(\mathbf{r},t) + \chi^{(3)}(\mathbf{r})E^{3}(\mathbf{r},t) + \dots$



$P(t) = \chi^{(1)}E(t) + \chi^{(2)}E^2(t) + \chi^{(3)}E^3(t) + \dots$



$P_i(t) = \chi_{ij}^{(1)} E_j(t) + \chi_{ijk}^{(2)} E_j(t) E_k(t) + \chi_{ijkl}^{(3)} E_j(t) E_k(t) E_l(t) + \dots$

$P_{i}(t) = \chi_{ij}^{(1)} E_{j}(t) + \chi_{ijk}^{(2)} E_{j}(t) E_{k}(t) + \chi_{ijkl}^{(3)} E_{j}(t) E_{k}(t) E_{l}(t) + \dots$

implicit sum over repeated Cartesian components $P_{i}(t) = \chi_{ij}^{(1)}E_{j}(t) + \chi_{ijk}^{(2)}E_{j}(t)E_{k}(t) + \chi_{ijkl}^{(3)}E_{j}(t)E_{k}(t)E_{l}(t) + \dots$ $vanishes \ if \ medium \ has centre-of-inversion symmetry$

$$E_i(t) = \sum_{\omega} E_i(\omega) e^{-i\omega t}$$

$$P_i(t) = \chi_{ij}^{(1)} E_j(t) + \chi_{ijk}^{(2)} E_j(t) E_k(t) + \chi_{ijkl}^{(3)} E_j(t) E_k(t) E_l(t) + \dots$$

$$E_i(t) = \sum_{\omega} E_i(\omega) e^{-i\omega t}$$

$$P_{i}(t) = \sum_{\omega} \chi_{ij}^{(1)}(-\omega;\omega)E_{j}(\omega)e^{-i\omega t}$$

$$+ \sum_{\omega,\omega'} \chi_{ijk}^{(2)}(-\omega - \omega';\omega,\omega')E_{j}(\omega)E_{k}(\omega')e^{-i(\omega+\omega')t}$$

$$+ \sum_{\omega,\omega',\omega''} \chi_{ijkl}^{(3)}(-\omega - \omega' - \omega'';\omega,\omega',\omega'')E_{j}(\omega)E_{k}(\omega')E_{l}(\omega'')e^{-i(\omega+\omega'+\omega'')t}$$

$$+ \dots$$

$$\begin{split} \chi_{ijk}^{(2)}(-\omega-\omega';\omega,\omega') \\ P_i(t) &= \sum_{\omega} \chi_{ij}^{(1)}(-\omega;\omega) E_j(\omega) e^{-i\omega t} \\ &+ \sum_{\omega,\omega'} \chi_{ijk}^{(2)}(-\omega-\omega';\omega,\omega') E_j(\omega) E_k(\omega') e^{-i(\omega+\omega')t} \\ &+ \sum_{\omega,\omega',\omega''} \chi_{ijkl}^{(3)}(-\omega-\omega'-\omega'';\omega,\omega',\omega'') E_j(\omega) E_k(\omega') E_l(\omega'') e^{-i(\omega+\omega'+\omega'')t} \\ &+ \dots \end{split}$$

$$\chi_{ijk}^{(2)}(-\omega - \omega'; \omega, \omega')$$

$$or \quad \chi_{ijk}^{(2)}(\omega, \omega'; \omega + \omega')$$

$$P_{i}(t) = \sum_{\omega} \chi_{ij}^{(1)}(-\omega; \omega) E_{j}(\omega) e^{-i\omega t}$$

$$+ \sum_{\omega, \omega'} \chi_{ijk}^{(2)}(-\omega - \omega'; \omega, \omega') E_{j}(\omega) E_{k}(\omega') e^{-i(\omega + \omega')t}$$

$$+ \sum_{\omega, \omega', \omega''} \chi_{ijkl}^{(3)}(-\omega - \omega' - \omega''; \omega, \omega', \omega'') E_{j}(\omega) E_{k}(\omega') E_{l}(\omega'') e^{-i(\omega + \omega' + \omega'')t}$$

$$+ \dots$$

$$\chi_{ijk}^{(2)}(-\omega - \omega'; \omega, \omega')$$

$$or \quad \chi_{ijk}^{(2)}(\omega, \omega'; \omega + \omega')$$

$$or \quad \chi_{ijk}^{(2)}(\omega, \omega')$$

$$P_{i}(t) = \sum_{\omega} \chi_{ij}^{(1)}(-\omega; \omega) E_{j}(\omega) e^{-i\omega t}$$

$$+ \sum_{\omega, \omega'} \chi_{ijk}^{(2)}(-\omega - \omega'; \omega, \omega') E_{j}(\omega) E_{k}(\omega') e^{-i(\omega + \omega')t}$$

$$+ \sum_{\omega, \omega', \omega''} \chi_{ijkl}^{(3)}(-\omega - \omega' - \omega''; \omega, \omega', \omega'') E_{j}(\omega) E_{k}(\omega') E_{l}(\omega'') e^{-i(\omega + \omega' + \omega'')t}$$

$$+ \dots$$

$$E_i(t) = \sum_{\omega} E_i(\omega) e^{-i\omega t}$$

$$P_{i}(t) = \sum_{\omega} \chi_{ij}^{(1)}(-\omega;\omega)E_{j}(\omega)e^{-i\omega t}$$

$$+ \sum_{\omega,\omega'} \chi_{ijk}^{(2)}(-\omega - \omega';\omega,\omega')E_{j}(\omega)E_{k}(\omega')e^{-i(\omega+\omega')t}$$

$$+ \sum_{\omega,\omega',\omega''} \chi_{ijkl}^{(3)}(-\omega - \omega' - \omega'';\omega,\omega',\omega'')E_{j}(\omega)E_{k}(\omega')E_{l}(\omega'')e^{-i(\omega+\omega'+\omega'')t}$$

$$+ \dots$$

$$E_i(t) = \sum_{\omega} E_i(\omega) e^{-i\omega t}$$

$$P_{i}(t) = \sum_{\omega} \chi_{ij}^{(1)}(\omega)E_{j}(\omega)e^{-i\omega t} + \sum_{\omega,\omega'} \chi_{ijk}^{(2)}(\omega,\omega')E_{j}(\omega)E_{k}(\omega')e^{-i(\omega+\omega')t} + \sum_{\omega,\omega',\omega''} \chi_{ijkl}^{(3)}(\omega,\omega',\omega'')E_{j}(\omega)E_{k}(\omega')E_{l}(\omega'')e^{-i(\omega+\omega'+\omega'')t} + \dots$$

$$E_i(t) = \sum_{\omega} E_i(\omega) e^{-i\omega t}$$

$$\begin{split} P_{i}(t) &= \sum_{\omega} \chi_{ij}^{(1)}(\omega) E_{j}(\omega) e^{-i\omega t} \\ &+ \sum_{\omega,\omega'} \chi_{ijk}^{(2)}(\omega,\omega') E_{j}(\omega) E_{k}(\omega') e^{-i(\omega+\omega')t} \\ &+ \sum_{\omega,\omega',\omega''} \chi_{ijkl}^{(3)}(\omega,\omega',\omega'') E_{j}(\omega) E_{k}(\omega') E_{l}(\omega'') e^{-i(\omega+\omega'+\omega'')t} \\ &+ \dots \end{split}$$

How calculate?





simple molecular models



simple molecular models



Calculate J(t) from minimal coupling Hamiltonian

$$\boldsymbol{J}(t) = rac{d\boldsymbol{P}(t)}{dt} \iff \boldsymbol{J}(\omega) = -i\omega\boldsymbol{P}(\omega)$$



simple molecular models



$$\boldsymbol{J}(t) = rac{d\boldsymbol{P}(t)}{dt} \iff \boldsymbol{J}(\omega) = -i\omega\boldsymbol{P}(\omega)$$

Adams and Blount



$$\begin{split} \boldsymbol{P} &= e \int \frac{d\boldsymbol{k}d\boldsymbol{k}'}{V} \sum_{n,m} a_n^{\dagger}(\boldsymbol{k}) \langle n\boldsymbol{k} | \boldsymbol{r} | m\boldsymbol{k'} \rangle a_m(\boldsymbol{k'}) \\ \langle n\boldsymbol{k} | \boldsymbol{r} | m\boldsymbol{k'} \rangle &= \delta(\boldsymbol{k} - \boldsymbol{k'}) \xi_{nm}(\boldsymbol{k}) + i \delta_{nm} \frac{\partial}{\partial \boldsymbol{k}} \delta(\boldsymbol{k} - \boldsymbol{k'}) \end{split}$$



Susceptibilities $\chi^{(2)}$ effects $\chi^{(3)}$ effects Quantum nonlinear optics Nonlinear optics and electronics Forbidden processes

Second harmonic generation

$$P_i(t) = \chi_{ijk}^{(2)}(-2\omega;\omega,\omega)E_j(\omega)E_k(\omega)e^{-2i\omega t} + c.c.$$

$$P_{i}(t) = \sum_{\omega} \chi_{ij}^{(1)}(-\omega;\omega)E_{j}(\omega)e^{-i\omega t}$$

$$+ \sum_{\omega,\omega'} \chi_{ijk}^{(2)}(-\omega - \omega';\omega,\omega')E_{j}(\omega)E_{k}(\omega')e^{-i(\omega+\omega')t}$$

$$+ \sum_{\omega,\omega',\omega''} \chi_{ijk}^{(3)}(-\omega - \omega' - \omega'';\omega,\omega',\omega'')E_{j}(\omega)E_{k}(\omega')E_{l}(\omega'')e^{-i(\omega+\omega'+\omega'')t}$$

$$+ \dots$$



















Second harmonic generation

GENERATION OF OPTICAL HARMONICS*

P. A. Franken, A. E. Hill, C. W. Peters, and G. Weinreich

The Harrison M. Randall Laboratory of Physics, The University of Michigan, Ann Arbor, Michigan (Received July 21, 1961)

The development of pulsed ruby optical masers^{1,2} has made possible the production of monochromatic (6943 A) light beams which, when focussed, exhibit electric fields of the order of 10^5 volts/cm. The possibility of exploiting this extraordinary intensity for the production of optical harmonics from suitable nonlinear materials is most appealing. In this Letter we present a brief discussion of the requisite analysis and a description of experiments in which we have observed the second harmonic (at ~ 3472 A) produced upon projection of an intense beam of 6943A light through crystalline quartz.

A suitable material for the production of optical harmonics must have a nonlinear dielectric coefficient and be transparent to both the fundamental optical frequency and the desired overtones. Since all dielectrics are nonlinear in high enough fields, this suggests the feasibility of utilizing materials such as quartz and glass. The dependence of polarization of a dielectric upon electric field E may be expressed schematically by

$$P = \chi E \left(1 + \frac{E}{E_1} + \frac{E^2}{E_2^2} + \cdots \right), \tag{1}$$

where E_1 , E_2 ... are of the order of magnitude of atomic electric fields (~10⁸ esu). If E is sinusoidal in time, the presence in Eq. (1) of terms of quadratic or higher degree will result in P con-

Table I. The square of the total p perpendicular to the direction of propagation of light through crystalline quartz.

Direction of incident beam	The square of the total p perpendicular to direction of propagation
$x (E_x = 0)$ $y (E_y = 0)$ $z (E_z = 0)$	$p_{y}^{2} + p_{z}^{2} = 0$ $p_{z}^{2} + p_{x}^{2} = \alpha^{2}E_{x}^{4}$ $p_{x}^{2} + p_{y}^{2} = \alpha^{2}(E_{x}^{2} + E_{y}^{2})^{2}$

(z is the threefold, or optic, axis; x a twofold axis). If a light beam traverses quartz in one of the three principal directions, Eqs. (2) predict the results summarized in Table I. The secondharmonic light should be absent in the first case, dependent upon incident polarization in the second case, and independent of this polarization in the third.

If an intense beam of monochromatic light is focussed into a region of volume V, there should occur an intensity I of second harmonic given (in Gaussian units) by

$$I \sim (\omega^4/c^3) (pv)^2 (V/v),$$
 (3)



FIG. 1. A direct reproduction of the first plate in which there was an indication of second harmonic. The wavelength scale is in units of 100 A. The arrow at 3472 A indicates the small but dense image produced by the second harmonic. The image of the primary beam at 6943 A is very large due to halation.




~ ^ ^ 2ω Z ω

 $P(2\omega) = \chi^{(2)} E(\omega) E(\omega)$

 2ω \sim \sim \sim \sim \sim

 $P(z, 2\omega) = \chi^{(2)} \left(E(\omega) e^{ik(\omega)z} \right) \left(E(\omega) e^{ik(\omega)z} \right)$

$$\begin{array}{c}
2\omega \\
\omega \\
\omega
\end{array}$$

$$P(z, 2\omega) = \chi^{(2)} \left(E(\omega) e^{ik(\omega)z} \right) \left(E(\omega) e^{ik(\omega)z} \right)$$
$$k(\omega) = \frac{\omega}{c} n(\omega)$$



$$P(z, 2\omega) = \chi^{(2)} E(\omega) E(\omega) e^{2ik(\omega)z}$$
$$k(\omega) = \frac{\omega}{c} n(\omega)$$



$$P(z, 2\omega) = \chi^{(2)} E(\omega) E(\omega) e^{2ik(\omega)z}$$
$$k(\omega) = \frac{\omega}{c} n(\omega) \qquad k(2\omega) = \frac{2\omega}{c} n(2\omega)$$



$$P(z, 2\omega) = \chi^{(2)} E(\omega) E(\omega) e^{2ik(\omega)z}$$
$$k(\omega) = \frac{\omega}{c} n(\omega) \qquad k(2\omega) = \frac{2\omega}{c} n(2\omega)$$

for constructive interference require: $k(2\omega) - 2k(\omega) = \frac{2\omega}{c}n(2\omega) - 2\frac{\omega}{c}n(\omega) = 0$ "phase matching condition"

$$E(t) = \left(E(\omega_S)e^{-i\omega_S t} + E(-\omega_S)e^{i\omega_S t}\right) + \left(E(\omega_I)e^{-i\omega_I t} + E(-\omega_I)e^{i\omega_I t}\right)$$

$$E(t) = \left(E(\boldsymbol{\omega}_{S})e^{-i\boldsymbol{\omega}_{S}t} + E(-\boldsymbol{\omega}_{S})e^{i\boldsymbol{\omega}_{S}t}\right) + \left(E(\boldsymbol{\omega}_{I})e^{-i\boldsymbol{\omega}_{I}t} + E(-\boldsymbol{\omega}_{I})e^{i\boldsymbol{\omega}_{I}t}\right)$$

second harmonic generation



 $P(2\omega_I) = \chi^{(2)} E(\omega_I) E(\omega_I)$

$$E(t) = \left(E(\boldsymbol{\omega}_{S})e^{-i\boldsymbol{\omega}_{S}t} + E(-\boldsymbol{\omega}_{S})e^{i\boldsymbol{\omega}_{S}t}\right) + \left(E(\boldsymbol{\omega}_{I})e^{-i\boldsymbol{\omega}_{I}t} + E(-\boldsymbol{\omega}_{I})e^{i\boldsymbol{\omega}_{I}t}\right)$$

second harmonic generation



.... and similarly for ω_s

$$E(t) = \left(E(\omega_S)e^{-i\omega_S t} + E(-\omega_S)e^{i\omega_S t}\right) + \left(E(\omega_I)e^{-i\omega_I t} + E(-\omega_I)e^{i\omega_I t}\right)$$



 $P(\omega_I + \omega_S) = \chi^{(2)} E(\omega_I) E(\omega_S)$

$$E(t) = \left(E(\omega_S)e^{-i\omega_S t} + E(-\omega_S)e^{i\omega_S t}\right) + \left(E(\omega_P)e^{-i\omega_P t} + E(-\omega_P)e^{i\omega_P t}\right)$$

difference frequency generation



 $P(\omega_P - \omega_S) = \chi^{(2)} E(\omega_P) E(-\omega_S)$



 $\chi^{(2)}(-2\omega;\omega,\omega)$







 $\chi^{(2)}(-2\omega;\omega,\omega)$

undepleted pump approximation



$$\chi_{ijk}^{(2)}(-\omega_1-\omega_2;\omega_1,\omega_2)$$



$$\chi_{ijk}^{(2)}(-\omega_1 - \omega_2; \omega_1, \omega_2)$$

second harmonic generation

$$\chi_{ijk}^{(2)}(-2\omega;\omega,\omega)$$



$$\chi_{ijk}^{(2)}(-\omega_1 - \omega_2; \omega_1, \omega_2)$$

second harmonic generation

 $\chi^{(2)}_{ijk}(-2\omega;\omega,\omega)$

linear electro-optic effect (Pockels effect)

 $\chi^{(2)}_{ijk}(-\omega;\omega,0)$

 $P(\omega) \propto E(\omega)E_{DC}$



$$\chi_{ijk}^{(2)}(-\omega_1-\omega_2;\omega_1,\omega_2)$$

second harmonic generation

 $\chi^{(2)}_{ijk}(-2\omega;\omega,\omega)$

linear electro-optic effect (Pockels effect)

 $\chi^{(2)}_{ijk}(-\omega;\omega,0)$

 $P(\omega) \propto E(\omega)E_{DC}$

difference frequency generation

$$\chi_{ijk}^{(2)}(-\omega_1+\omega_2;\omega_1,-\omega_2)$$



$$\chi_{ijk}^{(2)}(-\omega_1-\omega_2;\omega_1,\omega_2)$$

second harmonic generation

 $\chi^{(2)}_{ijk}(-2\omega;\omega,\omega)$

 $\chi_{iik}^{(2)}(-\omega;\omega,0)$

linear electro-optic effect (Pockels effect)

 $P(\omega) \propto E(\omega)E_{DC}$

 $P_{DC} \propto E(\omega)E(-\omega)$

difference frequency generation

 $\chi_{ijk}^{(2)}(-\omega_1+\omega_2;\omega_1,-\omega_2)$

optical rectification

 $\chi_{iik}^{(2)}(0;\omega,-\omega)$



$$\chi_{ijk}^{(2)}(-\omega_1-\omega_2;\omega_1,\omega_2)$$

second harmonic generation

 $\chi^{(2)}_{ijk}(-2\omega;\omega,\omega)$

 $\chi_{iik}^{(2)}(-\omega;\omega,0)$

linear electro-optic effect (Pockels effect)

 $P(\omega) \propto E(\omega)E_{DC}$

 $\left|P_{DC} \propto \left|E(\omega)\right|^2$

difference frequency generation

 $\chi_{ijk}^{(2)}(-\omega_1+\omega_2;\omega_1,-\omega_2)$

optical rectification

 $\chi_{iik}^{(2)}(0;\omega,-\omega)$

Susceptibilities $\chi^{(2)}$ effects $\chi^{(3)}$ effects Quantum nonlinear optics Nonlinear optics and electronics Forbidden processes



 $P(2\omega) = \chi^{(2)} E(\omega) E(\omega)$





$$E(t) = \left(E(\omega_S) e^{-i\omega_S t} + E(-\omega_S) e^{i\omega_S t} \right) + \left(E(\omega_P) e^{-i\omega_P t} + E(-\omega_P) e^{i\omega_P t} \right)$$

$$E(t) = \left(E(\omega_{S})e^{-i\omega_{S}t} + E(-\omega_{S})e^{i\omega_{S}t} \right) + \left(E(\omega_{P})e^{-i\omega_{P}t} + E(-\omega_{P})e^{i\omega_{P}t} \right)$$

four wave mixing



$$P(2\omega_P - \omega_S) = \chi^{(3)} E(\omega_P) E(\omega_P) E(-\omega_S)$$

$$E(t) = \left(E(\omega_{S})e^{-i\omega_{S}t} + E(-\omega_{S})e^{i\omega_{S}t} \right) + \left(E(\omega_{P})e^{-i\omega_{P}t} + E(-\omega_{P})e^{i\omega_{P}t} \right)$$

four wave mixing



$$P(2\omega_P - \omega_S) = \chi^{(3)} E(\omega_P) E(\omega_P) E(-\omega_S)$$

and many others!



 $\chi^{(3)}$ effects

four wave mixing

 $\chi_{ijkl}^{(3)}(-2\omega_P+\omega_S;\omega_P,\omega_P,-\omega_S)$

 $\chi^{(3)}$ effects

four wave mixing

$$\chi_{ijkl}^{(3)}(-2\omega_P+\omega_S;\omega_P,\omega_P,-\omega_S)$$

including CARS (coherent anti-Stokes Raman scattering)

 $\chi^{(3)}$ effects

four wave mixing

 $\chi_{ijkl}^{(3)}(-2\omega_P+\omega_S;\omega_P,\omega_P,-\omega_S)$

including CARS (coherent anti-Stokes Raman scattering)

self-phase modulation

 $\chi^{(3)}_{ijkl}(-\omega;\omega,\omega,-\omega)$

 $P(\omega) \propto E(\omega) \left| E(\omega) \right|^2$

 $\chi^{(3)}$ effects

four wave mixing $\chi^{(3)}_{ijkl}(-2\omega_P + \omega_S; \omega_P, \omega_P, -\omega_S)$

self-phase modulation

 $\chi_{iikl}^{(3)}(-\omega;\omega,\omega,-\omega)$

cross-phase modulation

 $\chi_{ijkl}^{(3)}(-\omega_1;\omega_1,\omega_2,-\omega_2)$

including CARS (coherent anti-Stokes Raman scattering)

 $P(\omega) \propto E(\omega) |E(\omega)|^2$

 $P(\omega_1) \propto E(\omega_1) |E(\omega_2)|^2$

 $\chi^{(3)}$ effects

four wave mixing $\chi^{(3)}_{ijkl}(-2\omega_P + \omega_S; \omega_P, \omega_P, -\omega_S)$

including CARS (coherent anti-Stokes Raman scattering)

 $P(\omega) \propto E(\omega) |E(\omega)|^2$

self-phase modulation $\chi^{(3)}_{ijkl}(-\omega;\omega,\omega,-\omega)$

cross-phase modulation

 $\chi_{ijkl}^{(3)}(-\omega_1;\omega_1,\omega_2,-\omega_2)$

including SRS (stimulated Raman scattering)

 $\chi^{(3)}$ effects

four wave mixing $\chi^{(3)}_{ijkl}(-2\omega_P + \omega_S; \omega_P, \omega_P, -\omega_S)$

self-phase modulation $\chi^{(3)}_{ijkl}(-\omega;\omega,\omega,-\omega)$

cross-phase modulation

 $\chi_{ijkl}^{(3)}(-\omega_1;\omega_1,\omega_2,-\omega_2)$

electro-optic effect (Kerr effect)

 $\chi_{ijkl}^{(3)}(-\omega;\omega,0,0)$

including CARS (coherent anti-Stokes Raman scattering)

 $P(\omega) \propto E(\omega) \left| E(\omega) \right|^2$

including SRS (stimulated Raman scattering)


Susceptibilities $\chi^{(2)}$ effects $\chi^{(3)}$ effects Quantum nonlinear optics Nonlinear optics and electronics Forbidden processes

$\hbar\omega_3$ $\hbar\omega_1$ $\hbar\omega_2$



 $\epsilon^{zz}(\omega)$





 $\epsilon_{z^z}^{zz}(\omega)$



 $\epsilon^{zz}(\omega)$

Classical nonlinear optics



nonlinear polarizations, envelope functions, coupled mode equations,...

Quantum nonlinear optics



Classical nonlinear optics



Hamiltonians, raising and lowering operators, Wigner functions,... nonlinear polarizations, envelope functions, coupled mode equations,...



difference frequency generation







spontaneous parametric down conversion



difference frequency generation



spontaneous four wave mixing





spontaneous parametric down conversion



difference frequency generation



spontaneous four wave mixing





spontaneous parametric down conversion



spontaneous four wave mixing



difference frequency generation







Kwiat et al. PRL 75 4337 (1995)

$$|\psi\rangle = \frac{1}{\sqrt{2}} \left(|H\rangle_{A} |V\rangle_{B} + |V\rangle_{A} |H\rangle_{B} \right)$$
$$|\psi\rangle \neq |\gamma\rangle_{A} |\sigma\rangle_{B}$$



quantum information processing





quantum simulation /computing



quantum mechanics



Photonic Quantum Computing





"Beyond photon pairs: nonlinear quantum photonics in the high-gain regime" Advances in Optics and Photonics 14, 291 (2022) Susceptibilities $\chi^{(2)}$ effects $\chi^{(3)}$ effects Quantum nonlinear optics Nonlinear optics and electronics Forbidden processes Energy (eV)





 $\epsilon_2^{\mathrm{zz}}(\omega)$

Bliefgy [CV]



$$\chi^{(3)}_{ijkl}(-\omega;\omega,\omega,-\omega)$$

acquires an imaginary part

mergy [cv]



$$\chi^{(3)}_{ijkl}(-\omega;\omega,\omega,-\omega)$$
 acquires an

imaginary part

$$P(\omega) \propto E(\omega) |E(\omega)|^2$$



mergy [cv]



$$\chi^{(3)}_{ijkl}(-\omega;\omega,\omega,-\omega)$$
 acquires an

imaginary part

$$P(\omega) \propto E(\omega) |E(\omega)|^2$$





$$\chi^{(2)}_{ijk}(-\omega_{\Sigma};\omega_1,\omega_2) =$$

$$\overline{\chi}_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2) + \frac{\sigma_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2)}{(-i\omega_{\Sigma})} + \frac{\eta_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2)}{(-i\omega_{\Sigma})^2}$$

with

 $\omega_{\Sigma} = \omega_1 + \omega_2$





Suppose $\omega_1 = \omega > \omega_{gap}$ $\omega_2 = -\omega$



$$\chi_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2) =$$

$$\overline{\chi}_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2) + \frac{\sigma_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2)}{(-i\omega_{\Sigma})} + \frac{\eta_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2)}{(-i\omega_{\Sigma})^2}$$

with

 $\omega_{\Sigma} = \omega_1 + \omega_2$





optical rectification







$$J_i(\omega_{\Sigma}) = \sigma_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2)E_j(\omega_1)E_k(\omega_2)$$

CW limit

$$J_i = 2\sigma_{ijk}^{(2)}(0;\omega,-\omega)E_j(\omega)E_k(-\omega)$$

"shift current"
"photovoltaic effect"











 $\omega_{\Sigma} \to 0$

$$\frac{dJ_i}{dt}(\omega_{\Sigma}) = \eta_{ijk}^{(2)}(-\omega_{\Sigma};\omega_1,\omega_2)E_j(\omega_1)E_k(\omega_2)$$

CW limit

$$\frac{dJ_i}{dt} = 2\eta_{ijk}^{(2)}(0;\omega,-\omega)E_j(\omega)E_k(-\omega)$$

"injection current" "circular photocurrent" "photogalvanic effect"



amplitude for transition $\propto \mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(\omega)$

probability for transition $\propto |\mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(\omega)|^2$



amplitude for transition $\propto \mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(\omega)$

probability for transition $\propto |\mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(\omega)|^2$

$$\mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E} = p_{cv}^{x}(\mathbf{k}) E^{x}(\omega) + p_{cv}^{z}(\mathbf{k}) E^{z}(\omega)$$


 $\begin{array}{c} \text{amplitude for transition} \\ \propto \quad \mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(\omega) \end{array}$

probability for transition $\propto |\mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(\omega)|^2$

$$\mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E} = p_{cv}^{x}(\mathbf{k}) E^{x}(\omega) + p_{cv}^{z}(\mathbf{k}) E^{z}(\omega)$$



amplitude for transition $\propto \mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(\omega)$

probability for transition $\propto |\mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(\omega)|^{2}$

$$\mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E} = p_{cv}^{x}(\mathbf{k}) E^{x}(\omega) + p_{cv}^{z}(\mathbf{k}) E^{z}(\omega)$$

$$\left|p_{cv}^{x}(\mathbf{k})E^{x}(\omega)+p_{cv}^{z}(\mathbf{k})E^{z}(\omega)\right|^{2}$$

can show interference effects





direction of current from crystal axis depends on helicity of beam



$$\chi_{ijkl}^{(3)}(-\omega_{\Sigma};\omega_1,\omega_2,\omega_3) =$$

$$\overline{\chi}_{ijkl}^{(3)}(-\omega_{\Sigma};\omega_{1},\omega_{2},\omega_{3}) + \frac{\sigma_{ijkl}^{(3)}(-\omega_{\Sigma};\omega_{1},\omega_{2},\omega_{3})}{(-i\omega_{\Sigma})} + \frac{\eta_{ijkl}^{(3)}(-\omega_{\Sigma};\omega_{1},\omega_{2},\omega_{3})}{(-i\omega_{\Sigma})^{2}}$$

with $\omega_{\Sigma} = \omega_1 + \omega_2 + \omega_3$



with $\omega_{\Sigma} = \omega_1 + \omega_2 + \omega_3$

 $\omega_{\Sigma} \to 0$



 $\omega_{\Sigma} = \omega_1 + \omega_2 + \omega_3$

 $\omega_{\Sigma} \to 0$



$$A_{cv}^{(1)}(\mathbf{k}) \propto \mathbf{p}_{cv}(\mathbf{k}) \cdot \mathbf{E}(2\omega)$$

$$A_{cv}^{(2)}(\mathbf{k}) \propto \sum_{n} \frac{\left[\mathbf{p}_{cn}(\mathbf{k}) \cdot \mathbf{E}(\omega)\right] \left[\mathbf{p}_{nv}(\mathbf{k}) \cdot \mathbf{E}(\omega)\right]}{\left[\omega_{c}(\mathbf{k}) + \omega_{v}(\mathbf{k}) - 2\omega_{n}(\mathbf{k})\right]}$$









road



spin-polarized current



Susceptibilities $\chi^{(2)}$ effects $\chi^{(3)}$ effects Quantum nonlinear optics Nonlinear optics and electronics Forbidden processes Forbidden processes

Centrosymmetric medium

$$\chi_{ijk}^{(2)} = 0$$

Forbidden processes

Centrosymmetric medium

$$\chi_{ijk}^{(2)} = 0$$



Centrosymmetric medium

$$\chi_{ijk}^{(2)} = 0$$

$$P_{i}(\boldsymbol{r}, 2\omega) = \Gamma_{ijkl} E_{j}(\boldsymbol{r}, \omega) \frac{\partial}{\partial x_{k}} E_{k}(\boldsymbol{r}, \omega)$$

Forbidden processes (..and allowed)



$$P_{i}(\boldsymbol{r},2\omega)=\Gamma_{ijkl}E_{j}(\boldsymbol{r},\omega)\frac{\partial}{\partial x_{k}}E_{k}(\boldsymbol{r},\omega)$$

Forbidden processes (..and allowed)

$$P_{i}(\mathbf{r}, 2\omega) = \Delta_{ijk}^{(2)} E_{j}(\mathbf{r}, \omega) E_{k}(\mathbf{r}, \omega) \delta(z - 0^{+})$$

$$z = 0^{+}$$
Centrosymmetric medium
$$\chi_{ijk}^{(2)} = 0$$

$$P_{i}(\mathbf{r}, 2\omega) = \Gamma_{ijkl} E_{j}(\mathbf{r}, \omega) \frac{\partial}{\partial x_{k}} E_{k}(\mathbf{r}, \omega)$$

Susceptibilities $\chi^{(2)}$ effects $\chi^{(3)}$ effects Quantum nonlinear optics Nonlinear optics and electronics Forbidden processes

"Optical alchemy"



